

QUANTUM NOISES

Q. D. S.

S. Attal

I. Quant. Dyn. Semigroups

II. The Probabilistic Fock Space

III. Quantum Noises

IV. Recovering Classical Noises

V. Quantum Langevin Equations

I Quantum Dynamical Semigroups

$\mathcal{P}_{nr}(f) = \mathcal{L}_h^n(f)$ discrete semigroup of C.P. maps

$$\mathcal{L}_h(f) = \sum_i \Pi_i(h) f \Pi_i(h)^*$$

with $\sum_i \Pi_i(h)^* \Pi_i(h) = I$.

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$$\lim_{h \rightarrow 0} \frac{\mathcal{L}_h(f) - f}{h} = L(f)$$

$$\mathcal{L}_h(f) = f + hL(f) + o(h)$$

I Quantum Dynamical Semigroups

$\mathcal{P}_{\hbar}(\rho) = \mathcal{L}^{\hbar}(\rho)$ discrete semigroup of C.P. maps

$$\mathcal{L}_{\hbar}(\rho) = \sum_i \Pi_i(\hbar) \rho \Pi_i(\hbar)^*$$

$$\text{with } \sum_i \Pi_i(\hbar)^* \Pi_i(\hbar) = \mathbb{I}.$$

$$\lim_{\hbar \rightarrow 0} \frac{\mathcal{L}_{\hbar}(\rho) - \rho}{\hbar} = \mathcal{L}(\rho)$$

$$\mathcal{L}_{\hbar}(\rho) = \rho + \hbar \mathcal{L}(\rho) + o(\hbar)$$

$$\Rightarrow \begin{cases} \Pi_0(\hbar) = \mathbb{I} + \hbar \mathcal{L}_0 + o(\hbar) \\ \Pi_i(\hbar) = \sqrt{\hbar} \mathcal{L}_i + o(\sqrt{\hbar}) \end{cases}$$

The condition $\sum_i \Pi_i(\hbar)^* \Pi_i(\hbar) = \mathbb{I}$ then becomes

$$\mathcal{L}_0 = iH - \frac{1}{2} \sum_k \mathcal{L}_k^* \mathcal{L}_k. \quad (H \text{ self-adjoint})$$

We get

$$\mathcal{L}(\rho) = i[H, \rho] - \frac{1}{2} \sum_k (L_k^\dagger L_k \rho + \rho L_k^\dagger L_k - 2L_k \rho L_k^\dagger)$$

Lindblad generator.

Typical generator of $(\mathcal{P}_t)_{t \in \mathbb{R}^+}$ semigroup of quantum channels.

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For observables:

$$\mathcal{L}^*(X) = -i[H, X] - \frac{1}{2} \sum_k (L_k^* L_k X + X L_k^* L_k - 2L_k^* X L_k).$$

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$$\bar{\Phi} = \Gamma_s(L^2(\mathbb{R}^+; \mathbb{C}))$$

vacuum Ω

creation op. $a^*(f)$, $f \in L^2(\mathbb{R}^+)$

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$$\chi_t = a^*(\mathbb{1}_{[0,t]}) \Omega$$

$$\bar{\Phi}_{[t_1, t_2]} = \Gamma_0(L^2([t_1, t_2])) \subset \bar{\Phi}$$

$$\bar{\Phi} \sim \bar{\Phi}_{[0, t_1]} \otimes \bar{\Phi}_{[t_1, t_2]} \otimes \dots \otimes \bar{\Phi}_{[t_n, +\infty[}$$

$(\chi_t)_{t \in \mathbb{R}^+}$ is unique s.t. $\chi_t - \chi_s \in \bar{\Phi}_{[s, t]}$ $\forall s < t$.

The products

$$(X_{t_1+h_1} - X_{t_1}) \otimes \cdots \otimes (X_{t_n+h_n} - X_{t_n})$$

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$$F = c\Omega + \sum_{n=1}^{\infty} \int_{0 < t_1 < \cdots < t_n} f(t_1, \dots, t_n) dX_{t_1} \otimes \cdots \otimes dX_{t_n}.$$

$$\|F\|^2 = |c|^2 + \sum_{n=1}^{\infty} \int_{0 < t_1 < \cdots < t_n} |f(t_1, \dots, t_n)|^2 dt_1 \cdots dt_n.$$

$dX_t = dX_{t_1} \otimes \cdots \otimes dX_{t_n}$ form a continuous basis of $\bar{\mathcal{F}}$.

$$\bar{\mathcal{F}} = \left(\bigotimes_{t \in \mathbb{R}^+} \mathbb{C}_{(t)}^2 \right)$$

where $\mathbb{C}_{(t)}^2$ is generated by Ω_t and dX_t .

Quantum noise :

$$\int_0^t H_s dA_s$$

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→ composition \Rightarrow domain problems

→ non commutation $H_t dA_t \neq dA_t H_t$

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First approach:

$$H_t = H_t \otimes \mathcal{I} \text{ on } \mathcal{F}_{[0,t]} \otimes \mathcal{F}_{(t,+\infty)} \text{ (easy)}$$

$$A_t - A_s = \mathcal{I} \otimes (A_t - A_s) \otimes \mathcal{I} \text{ on } \mathcal{F}_{[0,s]} \otimes \mathcal{F}_{(s,t]} \otimes \mathcal{F}_{(t,+\infty)}$$

$$H_{t_i} (A_{t_{i+1}} - A_{t_i}) = H_{t_i} \otimes (A_{t_{i+1}} - A_{t_i})$$

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But condition on (A_t) : far from obvious.

	Ω	dX_t
$da_0^o(t)$	$dt \Omega$	0
$da_i^o(t)$	dX_t	0
$da_0^i(t)$	0	$dt \Omega$
$da_i^i(t)$	0	dX_t

$$a_j^t(t) = \int_0^t da_j^t(s)$$

$$a_j^t(t) f = \int_0^t da_j^t(s) \int_{\sigma} f(\sigma) dX_{\sigma}$$

$$= \int_0^t \int_{\sigma} f(\sigma) da_j^t(s) dX_{\sigma}$$

$$a_j^{\epsilon}(t) = \int_0^t da_j^{\epsilon}(s)$$

$$\begin{aligned} a_j^{\epsilon}(t) f &= \int_0^t da_j^{\epsilon}(s) \int_{\sigma} f(\sigma) dX_{\sigma} \\ &= \int_0^t \int_{\sigma} f(\sigma) da_j^{\epsilon}(s) dX_{\sigma} \end{aligned}$$

\Rightarrow explicit formulas, defining concrete operators

$$[a_0^0(t) f](\sigma) = \int_0^t f(\sigma) ds = t f(\sigma) = [t f](\sigma)$$

$$[a_1^0(t) f](\sigma) = \sum_{\substack{\lambda \in \sigma \\ \lambda \leq t}} f(\sigma \setminus \lambda)$$

$$[a_0^1(t) f](\sigma) = \int_0^t f(\sigma \cup \{s\}) ds$$

$$[a_1^1(t) f](\sigma) = \#(\sigma \cap [0, t]) f(\sigma).$$

$$T_t = \int_0^t H_s da_j^t(s). \text{ well-defined.}$$

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$$S_t = \int_0^t K_s da_\ell^k(s)$$

$$T_t S_t = \int_0^t T_s dS_s + \int_0^t dT_s S_s + \int_0^t dT_0 dS_s$$

$$= \int_0^t T_s K_s da_\ell^k(s) + \int_0^t H_s S_s da_j^i(s)$$

$$+ \int_0^t H_s K_s da_j^i(s) da_\ell^k(s)$$

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where $da_0^1 da_1^0 = da_0^0$

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all other products = 0.

Wiener space $(\Omega, \mathcal{F}, \mathbb{P})$, (W_t) .

$$\mathcal{L}^2(\Omega) \ni F = \mathbb{E}[F] + \sum_{n=1}^{\infty} \int_{0 < t_1 < \dots < t_n} f(t_1, \dots, t_n) dW_{t_1} \dots dW_{t_n}$$

$$\|F\|^2 = |\mathbb{E}[F]|^2 + \sum_{n=1}^{\infty} \int_{0 < t_1 < \dots < t_n} |f(t_1, \dots, t_n)|^2 dt_1 \dots dt_n$$

$$\mathcal{L}^2(\Omega) \sim \Phi$$

$$\sigma_{W_t} = a_i^0(t) + a_i^1(t)$$

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$$\mathcal{L}^2(\Omega) \sim \Phi$$

$$\sigma_{W_t} = a_0^0(t) + a_0^1(t)$$

$$\text{Same } X_t = N_t - t$$

$$\sigma_{X_t} = a_0^0(t) + a_0^1(t) + a_1^1(t).$$

$$0. \mathcal{H}_s \otimes \underline{\Phi}$$

$$dV_t = \sum_{i,j} L_j^i V_t da_j^i(t)$$

$$L_j^i \in \mathcal{B}(\underline{\Phi}).$$

$$dV_t = L_0^0 V_t dt + \sum_{(i,j) \neq (0,0)} L_j^i V_t da_j^i(t).$$

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Thm

V_t unitary iff

- $L_1^1 = S - I$, S unitary
- $L_0^1 = -L_1^{0*} S$
- $L_0^0 = iH - \frac{1}{2} L_1^{0*} L_1^0$.

TRm

$$\text{tr}_2 \left[V_T (p \otimes |\Omega\rangle\langle\Omega|) V_T^* \right] = \mathcal{P}_T(p)$$

$$\mathcal{P}_T = e^{+L}$$

$$L(p) = -i [H, p] - \frac{1}{2} (L_1^{0*} L_1^0 p + p L_1^{0*} L_1^0 - 2 L_1^0 p L_1^{0*})$$

TRM

$$\mathbb{E}_t \left[V_T (p \otimes |R| \langle R |) V_T^* \right] = \mathcal{P}_T(p)$$

$$\mathcal{P}_T = e^{+L}$$

$$L(p) = -i [H, p] - \frac{1}{2} (L_1^{0*} L_1^0 p + p L_1^{0*} L_1^0 - 2 L_1^0 p L_1^{0*})$$

$$dV_T = (iH - \frac{1}{2} L^2) V_T dt + L V_T dW_T \quad (L \text{ s. a.})$$